

Uniformly Chaotic Finite-Range Lattice Models

And the Characterisation of the Set of Ground States Thereof

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Thermodynamic Formalism and Chaoticity

Controlling Markers Distribution

Building an Appropriate Structure

- Turing Machines as Tilings

- The Robinson Tiling(s)

- Structure for Entropy Control

Forcing Complex Structures

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Gibbs Measures on Finite Spaces

- Ω a *finite* set of states.
- $E : \Omega \rightarrow \mathbb{R}^+$ an *energy* function.
- β the inverse temperature.

Theorem (Variational Principle)

The distribution $\mu_\beta(\omega) \propto \exp(-\beta E(\omega))$ is the only maximiser of $\mu \mapsto H(\mu) - \beta\mu(E)$, with $H(\mu) := \sum -\log_2(\mu(\omega))\mu(\omega)$ the entropy.

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We call μ_β a Gibbs measure.

- At high temperatures, as $\beta \rightarrow 0$, we converge to the uniform distribution $\mathcal{U}(\Omega)$, that maximises H .
- At low temperatures, as $\beta \rightarrow \infty$, we converge to the uniform distribution $\mathcal{U}(\Omega^*)$, that maximises H among measures of minimal energy, supported by $\Omega^* := \arg \min(E)$.

Invariant Gibbs Measures on Lattice Models

- $\Omega_{\mathcal{A}} := \mathcal{A}^{\mathbb{Z}^d}$ the phase space, with \mathcal{A} a finite alphabet.
- $\mathbb{Z}^d \curvearrowright^{\sigma} \Omega_{\mathcal{A}}$ the shift action, so that $\sigma^x(\omega)_y = \omega_{y-x}$ for any $x, y \in \mathbb{Z}^d$ and $\omega \in \Omega_{\mathcal{A}}$.
- $\mathcal{M}_{\sigma}(\Omega_{\mathcal{A}})$ the set of invariant measures on $\Omega_{\mathcal{A}}$, such that $\mu \circ \sigma^x = \mu$ for any $x \in \mathbb{Z}^d$.
- $\varphi : \Omega_{\mathcal{A}} \rightarrow \mathbb{R}^+$ a continuous potential, the contribution of $0 \in \mathbb{Z}^d$ to the energy.

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Definition (Pressure Function)

Define the pressure $p_{\mu}(\beta) := h(\mu) - \beta\mu(\varphi)$,

with $h(\mu) := \lim_{n \rightarrow \infty} \frac{1}{n^d} H(\mu_{\llbracket 0, n-1 \rrbracket^d})$ the entropy per site.

Let $\mathcal{G}_{\sigma}(\beta) := \arg \max_{\mu \in \mathcal{M}_{\sigma}} p_{\mu}(\beta)$ the set of Gibbs measures.

- φ has finite range if it is *locally constant*, if $\varphi(\omega)$ only depends on $\omega_{\llbracket -r, r \rrbracket^d}$.

Limit Behaviour for Ground States

- We call $(\mu_\beta \in \mathcal{G}_\sigma(\beta))_{\beta>0}$ a *cooling trajectory* of the model.
- Denote $\mathcal{G}_\sigma(\infty) := \text{Acc}_{\beta \rightarrow \infty} \mathcal{G}_\sigma(\beta)$ the set of *ground states*, of accumulation points of all the cooling trajectories.
- $\mathcal{G}_\sigma(\infty)$ is a connected compact set (for the weak-* topology).

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Lemma

Assume that $X := \{\omega \in \Omega_{\mathcal{A}}, \forall x \in \mathbb{Z}^d, \varphi \circ \sigma^x(\omega) = 0\} \neq \emptyset$.

Then $\mathcal{G}_\sigma(\infty) \subset \mathcal{M}_\sigma(X)$, and the ground states have maximal entropy h in $\mathcal{M}_\sigma(X)$.

- Measures that maximise h in $\mathcal{M}_\sigma(X)$ are not necessarily in $\mathcal{G}_\sigma(\infty)$.

What can we ask about $\mathcal{G}_\sigma(\infty)$?

Stability and Chaos

Definition (Stability)

A model is stable if all the cooling trajectories converge to the same limit.

Definition (Chaoticity)

A model is chaotic if there is no converging cooling trajectory.

Definition (Uniformity)

A model is uniform if all the cooling trajectories have the same accumulation set.

Recap of Behaviours

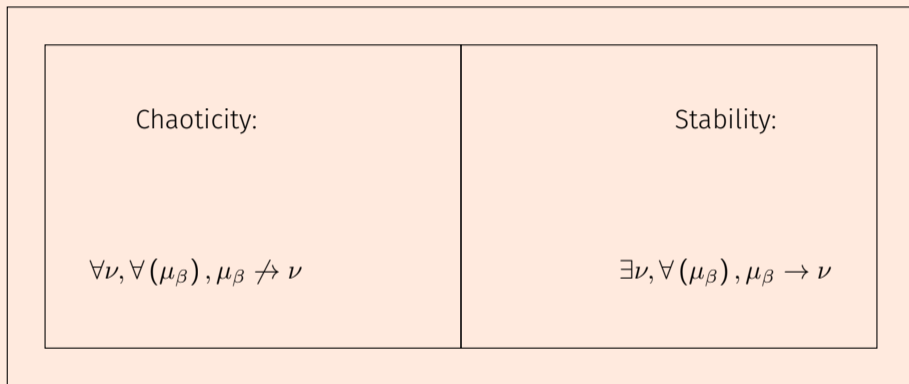


Figure 1: Inventory and comparison of model behaviours.

Recap of Behaviours

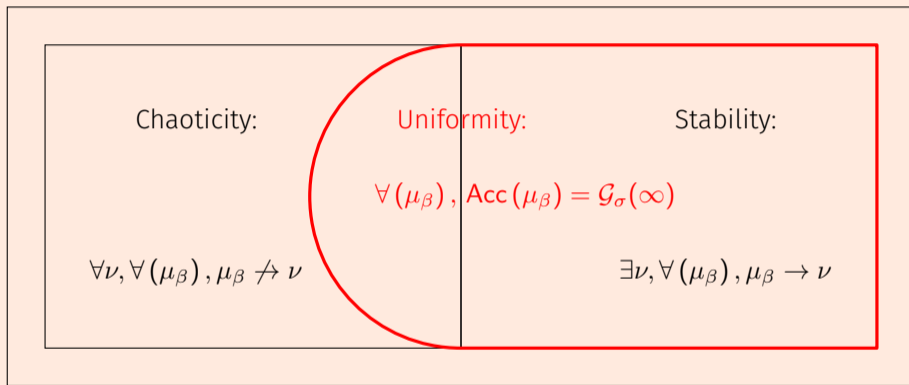


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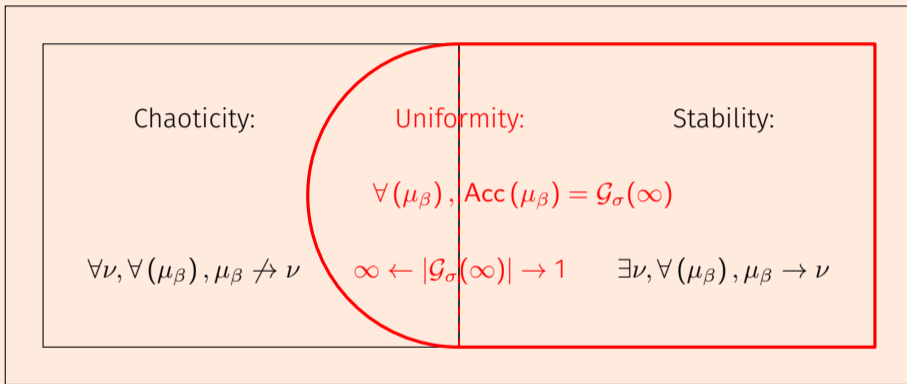


Figure 1: Inventory and comparison of model behaviours.

Current Knowledge

Lemma

A one-dimensional finite range model induces a stable model.

Theorem (Chazottes and Hochman 2010)

There exists a one-dimensional potential inducing a chaotic model.

There exists a three-dimensional finite range potential inducing a chaotic model.

Theorem (Chazottes and Shinoda 2020; Barbieri et al. 2022)

There exists a two-dimensional finite range potential inducing a chaotic model.

Realisation Result on the Limit Set

- Remind that $\mathcal{G}_\sigma(\infty)$ must be connected.
- When φ is a *computable* potential inducing a uniform model, $\mathcal{G}_\sigma(\infty)$ must be a Π_2 -computable set.

Theorem (Gayral, Sablik, and Taati 2023)

There exists a class of two-dimensional finite range potentials, inducing uniform models both stable and chaotic.

More precisely, we can realise any connected Π_2 -computable compact set X as $\mathcal{G}_\sigma(\infty)$, up to a fixed computable affine homeomorphism.

Controlling Markers Distribution

General Idea for Chaoticity

We have two measures $\mu \neq \mu'$ s.t. $d(\mu, \mu') \geq r$ and:

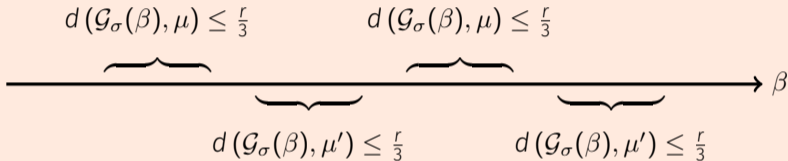


Figure 2: Alternating between mutually exclusive adherence values on non-overlapping intervals.

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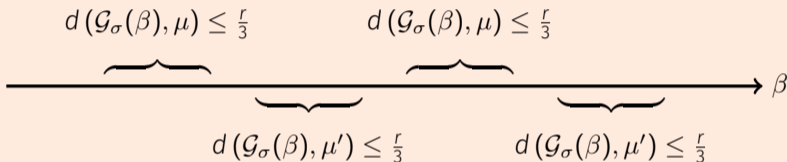


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Thus $\text{Acc}(\mu_\beta)$ intersects the disjoint neighbourhoods of both μ and μ' .

General Idea for Uniformity

We want (μ_n) and $\varepsilon_n \rightarrow 0$ s.t.:

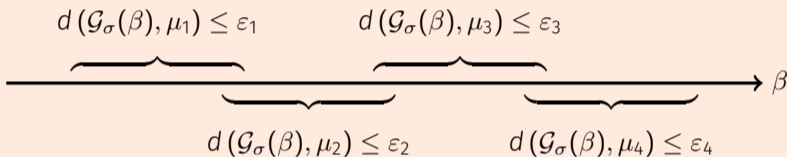


Figure 3: Contracting tube of measures with overlapping intervals.

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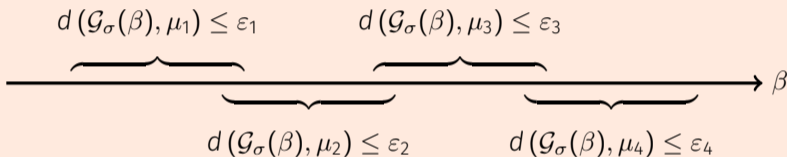


Figure 3: Contracting tube of measures with overlapping intervals.

Thus $\text{Acc}(\mu_\beta) = \mathcal{G}_\sigma(\infty) = \text{Acc}(\mu_n)$.

From Thermodynamics to Combinatorics

- \mathcal{F} a finite set of *forbidden patterns* $w \in \mathcal{A}^{l(w)}$, each on a finite window $l(w) \in \mathbb{Z}^d$.
- $p \in \mathcal{A}^l$ is locally admissible if no translation of a forbidden pattern occurs within it.
- \mathcal{F} induces a *subshift of finite type* (SFT) $X_{\mathcal{F}} \subset \Omega_{\mathcal{A}}$, closed and shift-invariant, made of configurations that are globally admissible.

Example

In one dimension, let $\mathcal{A} = \{0, 1\}$ and $\mathcal{F} = \{100, 101\}$. Then:

- $0^{\mathbb{Z}} \in X_{\mathcal{F}}, 1^{\mathbb{Z}} \in X_{\mathcal{F}}, \dots 000111 \dots \in X_{\mathcal{F}},$
- 10 is locally admissible, but doesn't occur in any $\omega \in X_{\mathcal{F}}$.

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Lemma

Assume that $X_{\mathcal{F}} \neq \emptyset$, and let $\varphi := \mathbb{1}_{\mathcal{F} \text{ covers } 0}$ the induced finite range potential.

Then $\mathcal{G}_{\sigma}(\infty) \subset \mathcal{M}_{\sigma}(X_{\mathcal{F}})$, and the ground states have maximal entropy h in $\mathcal{M}_{\sigma}(X_{\mathcal{F}})$.

Control of Markers on a Temperature Interval

Definition (Marker Set with Margin Factor τ)

A marker set $Q \subset \mathcal{A}^{I_\ell}$ (with $I_\ell := \llbracket 0, \ell - 1 \rrbracket^d$) is made of non-overlapping patterns, s.t. any locally admissible $\omega \in \mathcal{A}^{I_{(2+\tau)\ell-1}}$ must contain a marker somewhere.

Theorem (Adapted from Chazottes and Hochman 2010)

Denote G_n the locally admissible tilings of I_n , and μ_Q the cond. measure of μ on Q . We have constants C, C' s.t. for any marker set Q and $\varepsilon, \kappa > 0$, if

$$\frac{\log_2(\#G_n)}{\#I_n} \geq (1 - \kappa) \frac{\log_2(\#Q)}{\#I_\ell} \quad \text{and} \quad \beta \in \left[C \frac{\#I_\ell}{\varepsilon}, C' n \varepsilon \right]$$

then, for any $\mu \in \mathcal{G}_\sigma(\beta)$:

$$\mu(Q \text{ covers } 0) = 1 - O(\varepsilon + \tau) \quad \text{and} \quad H(\mu_Q) \geq (1 - 2\kappa) \log_2(\#Q) - H(\kappa) - O(\varepsilon + \tau).$$

Building an Appropriate Structure (*aka* LEGO for Grownups)

Turing Machines as Tilings

Turing Machines

Turing machines are a model equivalent to real-life computers and algorithms.

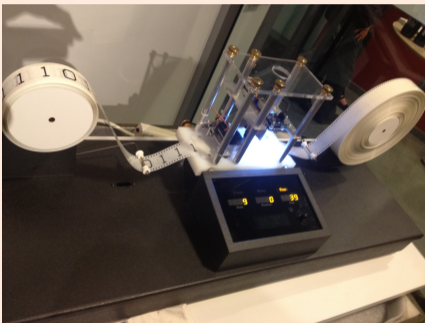


Figure 4: Real-life Turing machine
(Source: wikimedia.org)

Formally, M is made of:

- internal states Q ,
- an initial state $q_0 \in Q$,
- accepting states $Q_A \subset Q$,
- rejecting states $Q_R \subset Q$,
- an input alphabet \mathcal{A} ,
- a tape alphabet $\Gamma \supset \mathcal{A} \sqcup \{\#\}$,
- a transition function $\delta : Q \times \Gamma \rightarrow Q \times \Gamma \times \{L, R\}$.

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The Robinson Tiling(s)

Canonical Robinson Tiling (Non-Overlapping Markers)

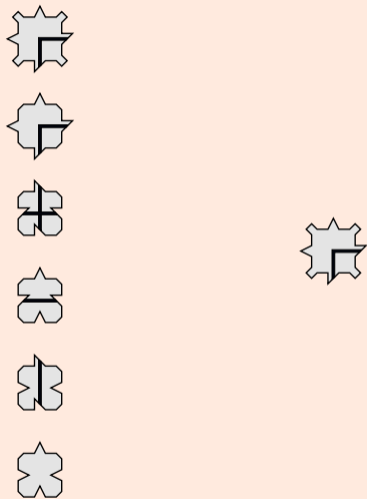


Figure 6: Hierarchical structure of the Robinson tiling.

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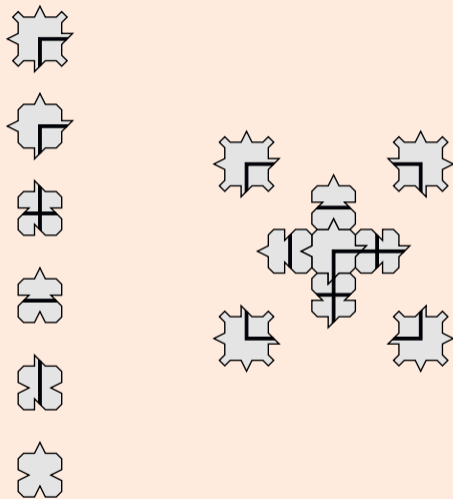


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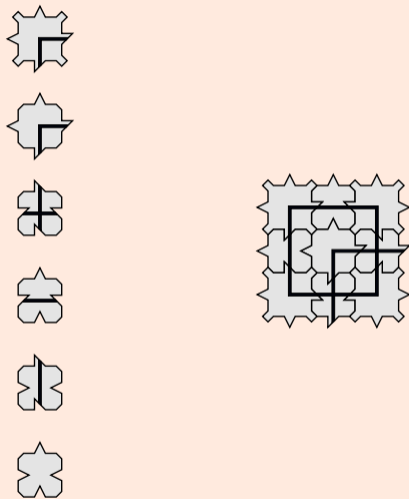


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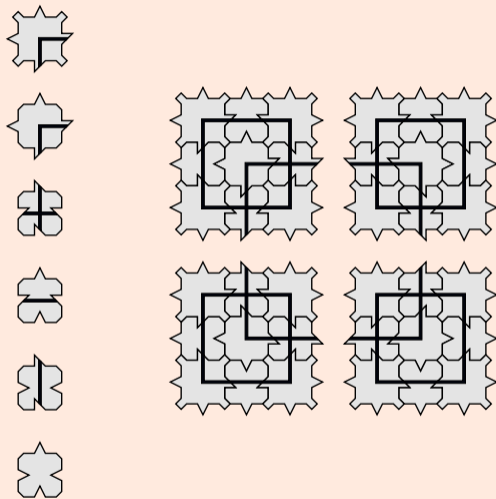


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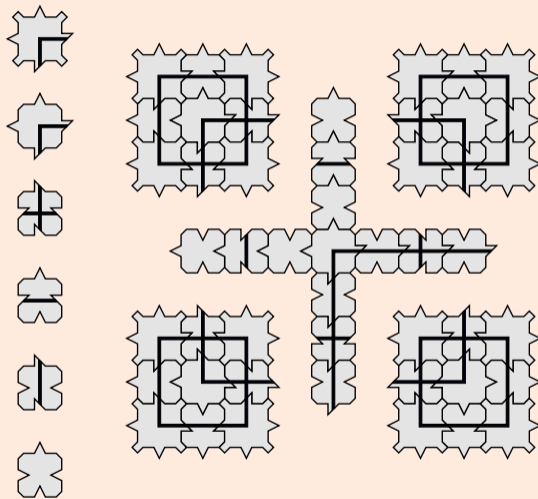


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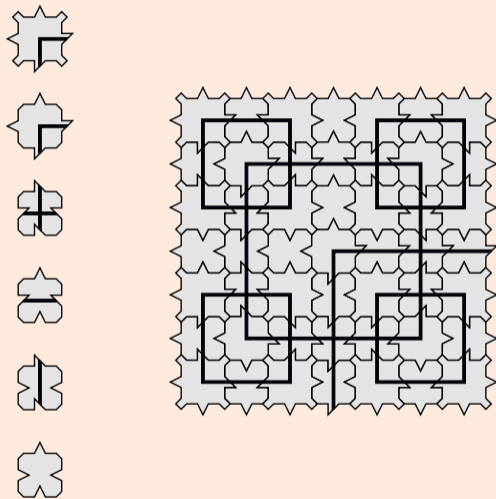


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Enhanced Robinson Tiling (Markers with Reconstruction)

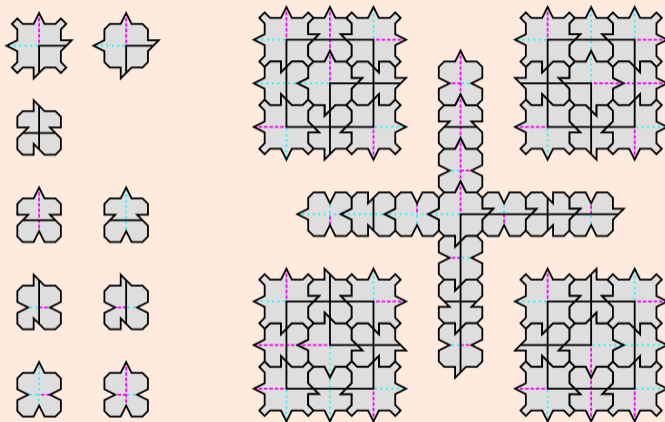


Figure 7: A Robinson variant, with strengthened local rules.

Two-Coloured Robinson for Turing Machines (Markers with Computation Area)

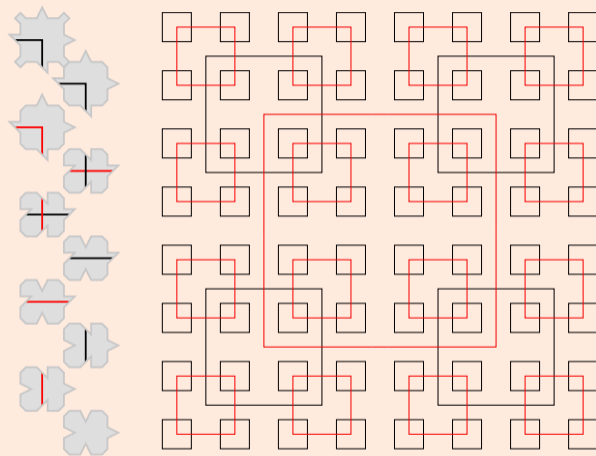


Figure 8: Alternating Red-Black structure,

Two-Coloured Robinson for Turing Machines (Markers with Computation Area)

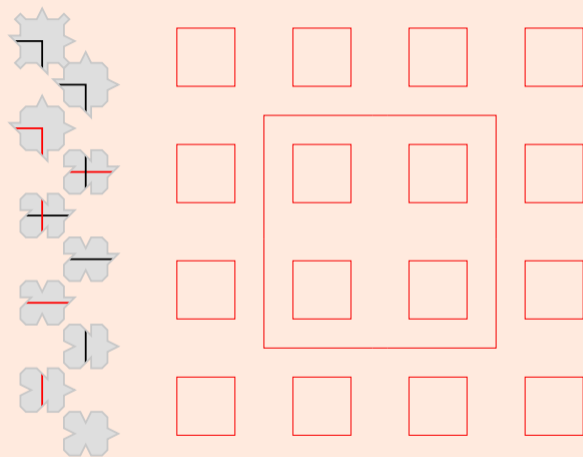


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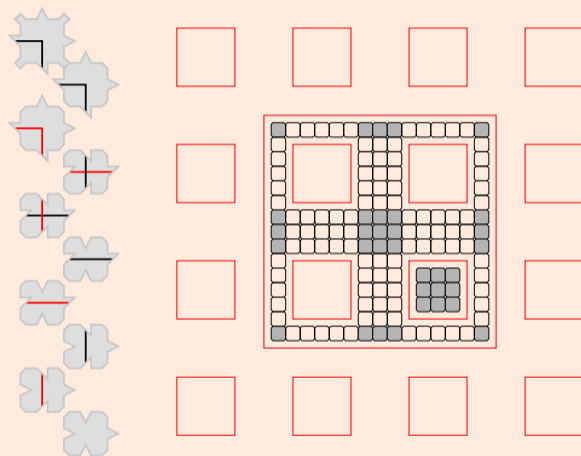


Figure 8: Alternating Red-Black structure, with a sparse computation area.

Structural Properties of the Base Layer

- The n -macro-tile has a length $\ell_n = 2^n - 1$.
- The n -macro-tiles are non-overlapping.
- Any locally admissible window of length $2\ell_n + 5$ contains a n -macro-tile. (Gayral, Sablik, and Taati 2023, Lemma 29)
- The N -th Red square occurs in a $(2N + 1)$ -macro-tile.
- The N -th Red square has a length $4^N + 1$.
- The N -th Red square has a sparse computing area of size $2^N + 1$.

Building an Appropriate Structure (*aka* LEGO for Grownups)

Structure for Entropy Control

Hot and Frozen Areas

Red squares may be Blocking, with a Hot exterior and Frozen core.
The rest must locally synchronise on Hot or Frozen.

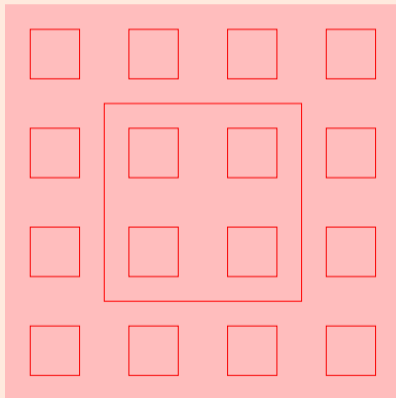


Figure 9: Admissible configurations for Hot and Frozen squares.

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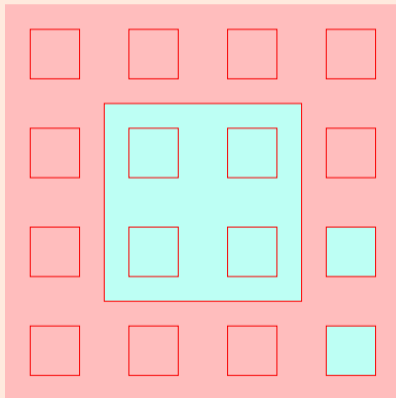


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Blockable Scales

We (can) unary encode N as an input for computations in the N -th Red square.
We check whether $N = 3^k$. If not, the Red square *cannot* be Blocking.

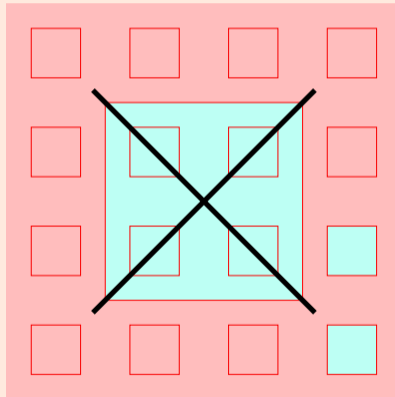


Figure 10: The 2nd scale of Red squares cannot be Blocking.

Scales for the Marker Sets

- Q_k the set of $n_k := (2 \times 3^k + 1)$ Robinson macro-tiles on the window $B_k := I_{\ell_{n_k}}$, the 3^k -th scale of locally admissible tiles with Red squares.
- A $(k + 1)$ -marker is a grid of $16^{3^k} \times 16^{3^k}$ smaller k -markers.
- This structure has positive entropy as each 0-marker, which occur periodically, can have a different state (either Hot or Blocking).

Odometer

We implement an odometer in k -markers, that cycles with period $t_k = 2^{\lfloor \log_2(\lfloor \log_2(k) \rfloor) \rfloor}$, so that Red squares are Blocking once for each cycle.

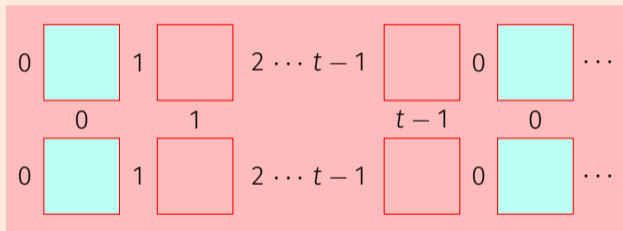


Figure 11: The repartition of Frozen squares is forced by the odometer.

The Red square of a $(k + 1)$ -marker initialises k -markers at 0 on one side.

Repartition of Frozen Tiles

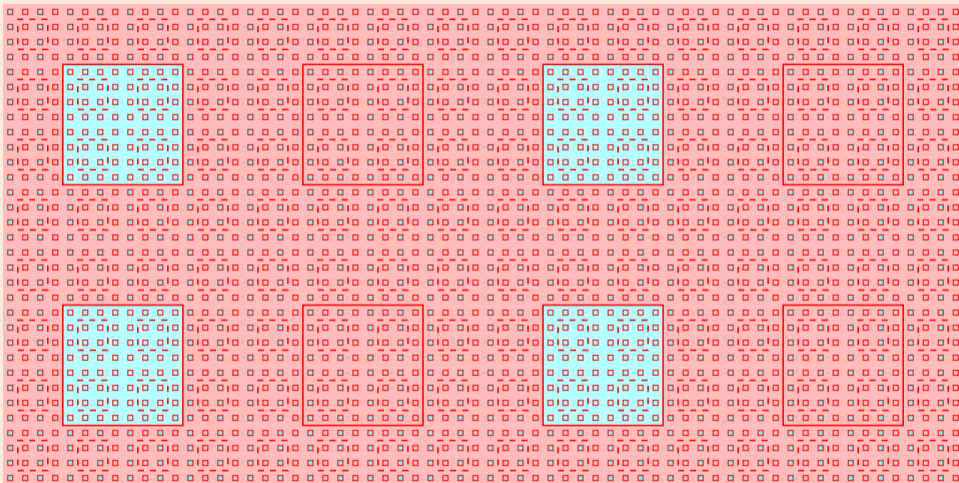


Figure 12: Approximation of a *Total Perspective Vortex*.
(One 2-marker would be a 4096×4096 grid of such 1-markers.)

Density of Frozen Tiles

The average scale of Blocking squares in a k -marker goes to ∞ as $k \rightarrow \infty$.

Proposition (Gayral, Sablik, and Taati 2023, Propositions 33 and 34)

Fix a microscopic scale m .

The proportion of non-Frozen m -markers in a k -marker is of order:

$$\prod_{j=m+1}^k \left(1 - \frac{1}{4t_j}\right) \xrightarrow{k \rightarrow \infty} 0.$$

Thus, generically, a globally admissible tiling is totally Frozen.

We are back to a uniquely ergodic zero-entropy case.

However, this rigid structure, with gaps in the scales, will allow us to slow down the speed of $\frac{\log_2(\#Q_k)}{\#B_k} \rightarrow 0$.

Words and Entropy

Encode a letter on Red lines so that:

- Blocking and Hot squares are labelled $\mathbf{0}$,
- Frozen squares are labelled $\pm\mathbf{1}$,
- Neighbouring Frozen squares synchronise their bit.

A Blocking k -marker central square encodes a binary word of length $3^k - 1$.

Generically, a (Frozen) tiling encodes a sequence of bits in $\{\pm\mathbf{1}\}^{\mathbb{N}}$.

Globally admissible tilings still have zero-entropy,
but now we have a source of entropy for locally admissible markers.

Counting Markers

Let $Q_k = Q_k^H \sqcup Q_k^B \sqcup Q_k^F$ depending on whether the Red square is Hot, Blocking or Frozen.

Proposition (Gayral, Sablik, and Taati 2023, Lemma 31, Propositions 42 and 43)

We have:

- $\#Q_k^H \approx C_k^{16^{3^k}}$ with $2^{4^{-k}} \leq C_k \leq 2$,
- $\#Q_k^B \approx (\#Q_k^H)^{\frac{3}{4}}$,
- $\#Q_k^F \leq C^{4^{3^k}}$ for some $C > 1$.

Thus, $\#Q_k \approx \#Q_k^H$.

A (Uniformly) Stable Structure

We conclude that μ_{Q_k} is close to the uniform distribution on Q_k^H .

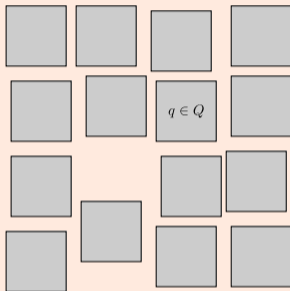


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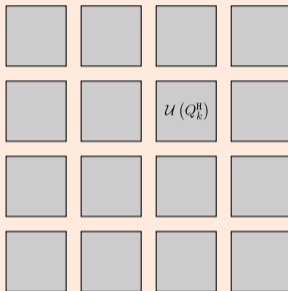


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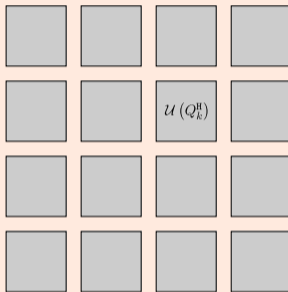


Figure 13: In the weak-* topology, Gibbs measures are approximately grids of uniform markers.

The induced model is uniform, stable, and the limit measure corresponds to $\mathcal{U}(\{\pm 1\}^{\mathbb{N}})$.

Forcing Complex Structures

Forcing a Distribution on Words

We can embed a Turing machine on a new layer to simulate a *non-uniform* distribution on the word encoded in each Blocking square.

This will easily give us uniformly chaotic models, e.g. by simulating δ_\emptyset , then δ_{11} , δ_{000} and so on, so that $\mathcal{G}_\sigma(\infty)$ corresponds to $[\delta_{\emptyset^{\mathbb{N}}}, \delta_{1^{\mathbb{N}}}]$.

What kind of kind of sets $\mathcal{G}_\sigma(\infty)$ we can obtain for this class of uniform models?

Computational Complexity of Uncountable Sets

Let (X, d) a metric space with a countable dense basis \mathcal{B} .

Definition

Let $Y \subset X$ be a closed set and $\mathcal{N}(Y) := \{(x, r) \in \mathcal{B} \times \mathbb{Q}^{+*}, \overline{B(x, r)} \cap Y \neq \emptyset\}$.

The set Y is said to be Π_k -computable *iff* the countable set $\mathcal{N}(Y)$ is, *i.e.* there is a computable φ such that:

$$(x, r) \in \mathcal{N}(Y) \Leftrightarrow \forall y_1, \exists y_2, \forall y_3, \dots, \varphi(x, r, y_1, \dots, y_k)$$

Here, for invariant measures $\mathcal{M}_\sigma(\Omega_A)$ with the weak-* topology, we use the periodic measures $\widehat{\delta}_w$, with $w \in \mathcal{A}^{\llbracket 0, n-1 \rrbracket^d}$, as a basis \mathcal{B} .

Uniform Upper Bound

Let φ a computable potential, inducing a uniform model.

Proposition (Gayral, Sablik, and Taati 2023, Proposition 3)

There is a sequence $\beta_k \rightarrow \infty$ such that $\text{diam}(\mathcal{G}_\sigma(\beta_k)) \rightarrow 0$ and $\mathcal{G}_\sigma(\infty) = \text{Acc}(\mathcal{G}_\sigma(\beta_k))$.

Without loss of generality, we have rational parameters $\beta_k \in \mathbb{Q}$.

Theorem (Gayral, Sablik, and Taati 2023, Theorem 17)

We have $\overline{B(x, r)} \cap \mathcal{G}_\sigma(\infty) \neq \emptyset$ iff:

$$\forall \varepsilon \in \mathbb{Q}^{+*}, \forall \beta_0 \in \mathbb{Q}^{+*}, \quad \exists \beta \in \mathbb{Q}_{>\beta_0}^{+*}, \exists y \in \mathcal{B}, \\ \mathcal{G}_\sigma(\beta) \subset B(y, \varepsilon) \text{ and } B(y, \varepsilon) \cap \overline{B(x, r)} \neq \emptyset.$$

Consequently, we have a Π_2 upper bound on the complexity of $\mathcal{G}_\sigma(\infty)$.

Equivalent Characterisation of Π_2





Proposition (Gayral, Sablik, and Taati 2023, Proposition 5)

There is a characterisation of Π_2 -computable sets through accumulation points:

$$\begin{aligned}
 Y \in \Pi_2 & \Leftrightarrow Y = \text{Acc}(x_n) \text{ with } (x_n) \in \mathcal{B}^{\mathbb{N}} \text{ computable.} \\
 Y \in \Pi_2 \text{ and connected} & \Leftrightarrow Y = \text{Acc}(x_n) \text{ with } (x_n) \in \mathcal{B}^{\mathbb{N}} \text{ computable,} \\
 & \text{and } d(x_n, x_{n+1}) \rightarrow 0.
 \end{aligned}$$

Thus, we can embed the Turing machine computing any such sequence, to obtain any Π_2 connected subset of $\mathcal{M}(\{\pm 1\}^{\mathbb{N}})$ encoded in $\mathcal{G}_\sigma(\infty)$.

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THE END OF PRESENTATION

ONE MORE SLIDE:

Thank you.